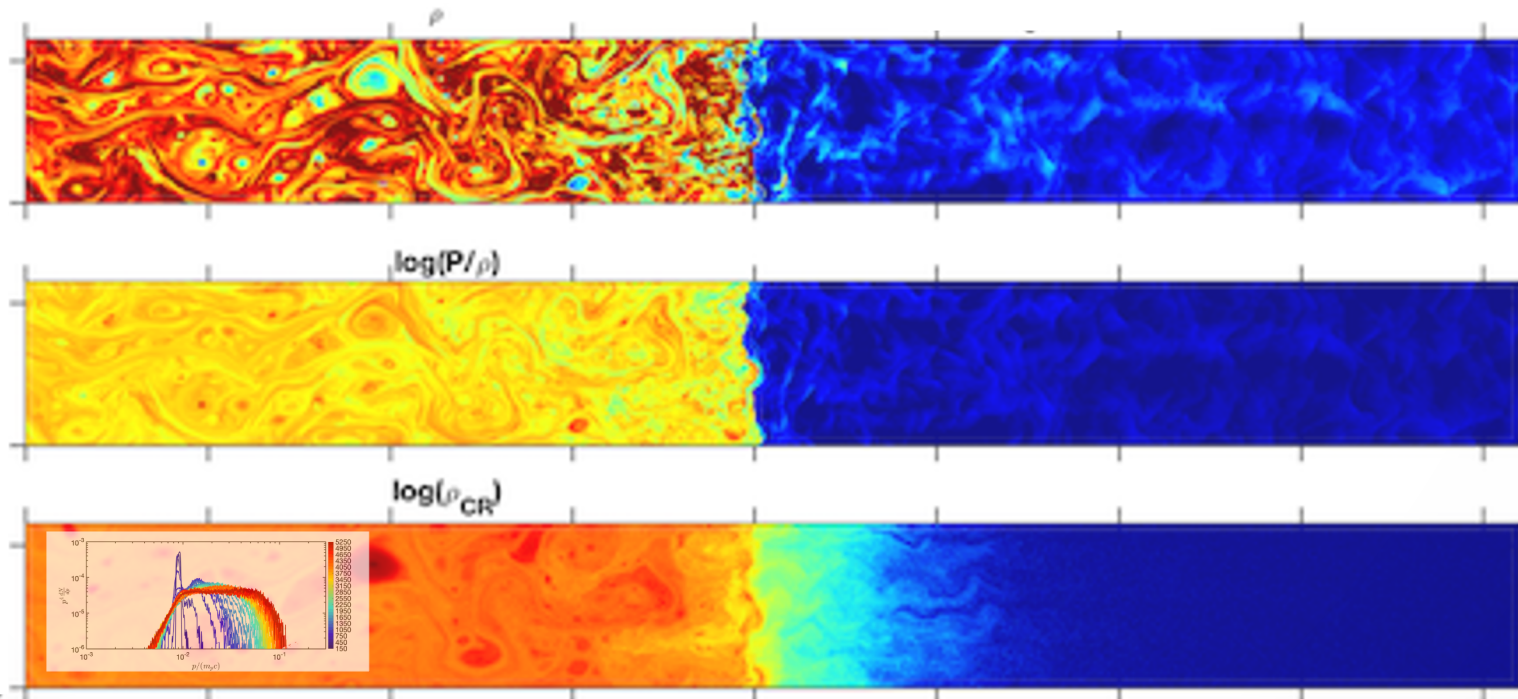


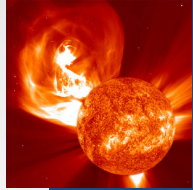
Micro- and multi-scale simulations of particle acceleration

Illya Plotnikov

Ateliers Astroplasma IRAP, 10 Feb 2021



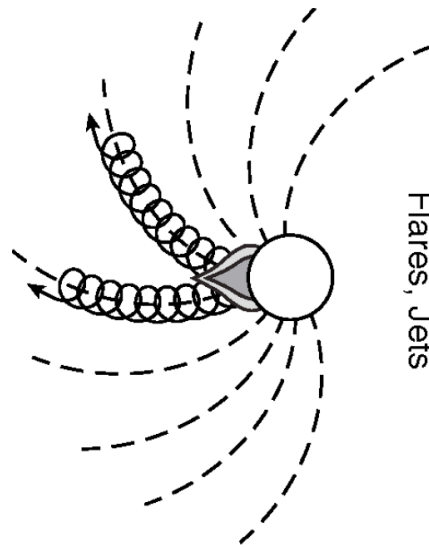
Examples of involved environments



Flavors of particle acceleration in space/astrophysics

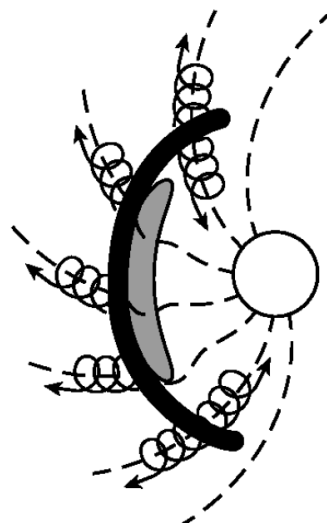
- Solar flares / CMEs

Measured in-situ as **impulsive** event
(acceleration in the small flare region)

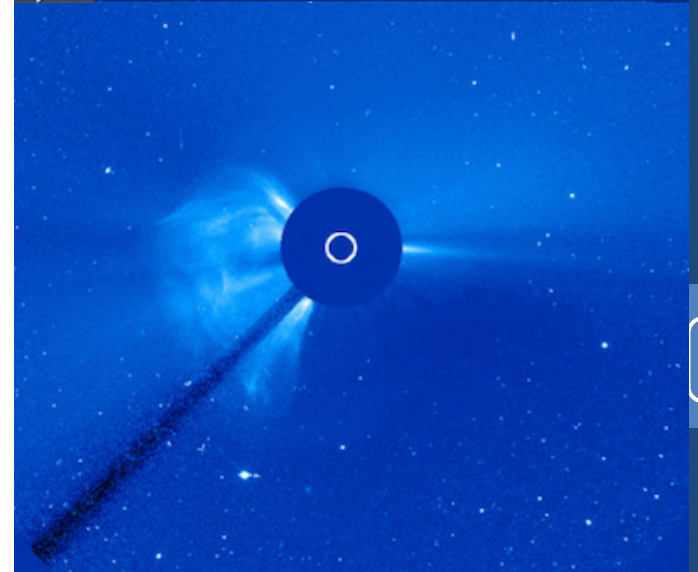
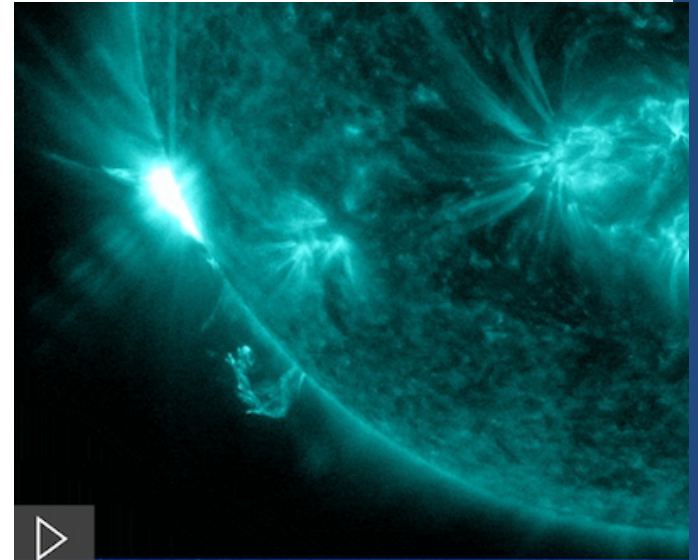


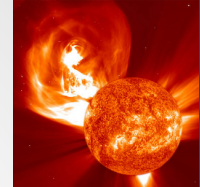
Flares, Jets

Measured in-situ as **gradual** event
(acceleration in the large CME region)



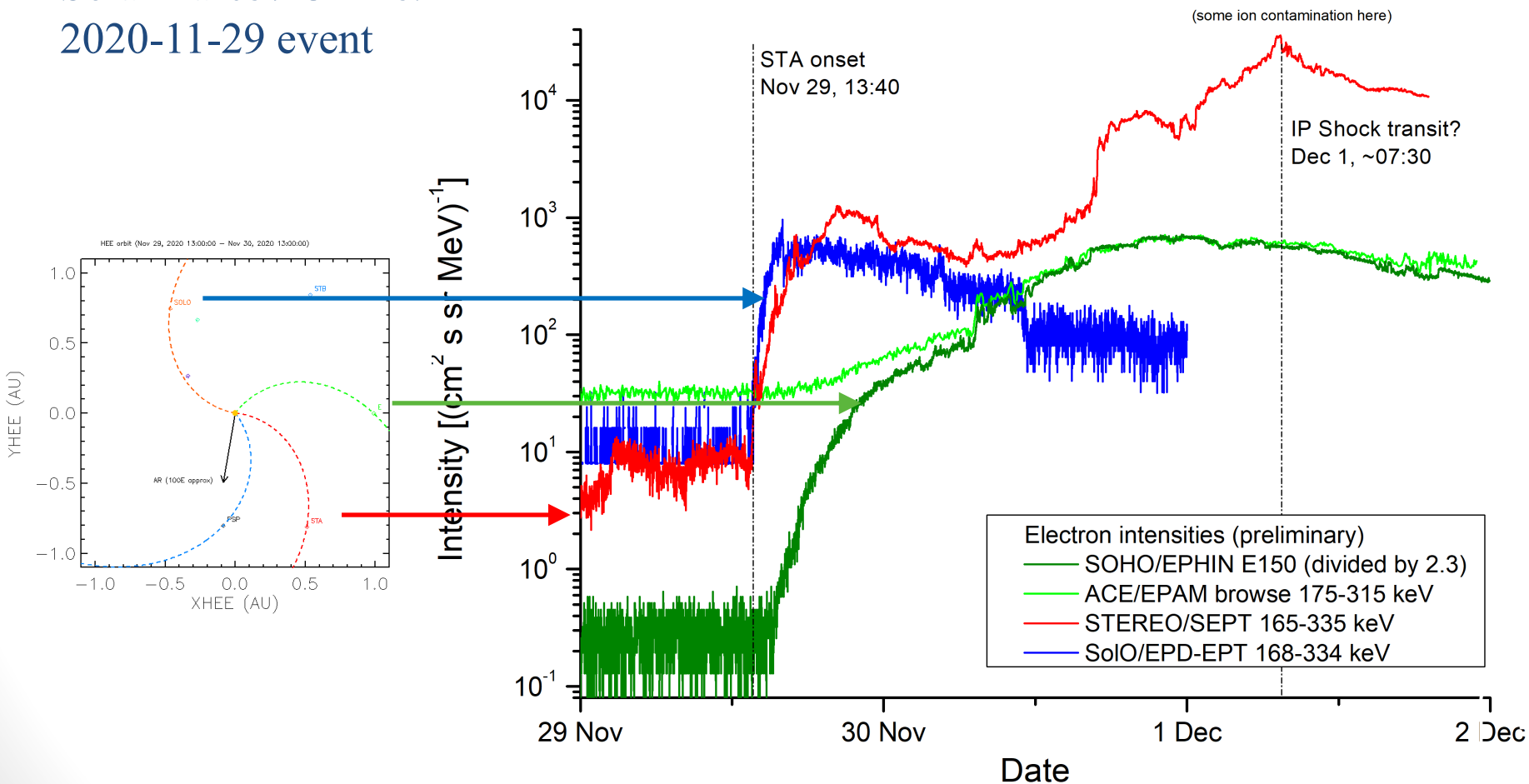
CME Shocks





Flavors of Fermi acceleration in space/astrophysics

- Solar flares / CMEs:
2020-11-29 event



Times delays : evolving magnetic connectivity of the expanding shock to remote probes
Max intensity: history of shock strength & local plasma density ?

Flavors of Fermi acceleration in space/astrophysics

Acceleration – FR II radio-galaxies



Faranoff-Riley II radio-galaxy Cygnus A

acceleration in the central AGN:
unipolar inductor, shocks in blazar zone...

AGN

hot spots

acceleration in hot spots:
mildly relativistic shocks...

relativistic
jets

acceleration in the jets: shocks, shear...
(... only FR II carry relativistic jets)

lobes

acceleration in the lobes:
no strong shock... stochastic Fermi acceleration

→ too few such sources in the GZK volume to account for the bulk of UHECRs...
unless they are heavy nuclei (← large magnetic deflection)...

e.g. Ptuskin et al. 13, Aloisio et al. 11

Flavors of Fermi acceleration in space/astrophysics

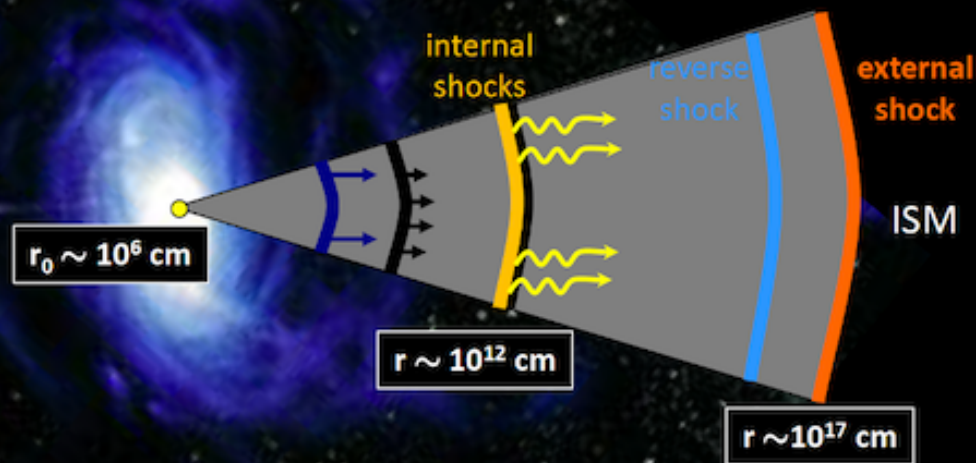
Gamma-ray bursts



... gamma-ray bursts: burst (<1 sec → 1000sec) of gamma radiation, with erratic time behavior in the MeV range, followed by a slowly decaying afterglow

... at the origin: collapse of massive stars (long?), coalescence of compact objects (short)?

... canonical description: narrow jet accelerated to large Lorentz factor $\Gamma \sim 100-1000$



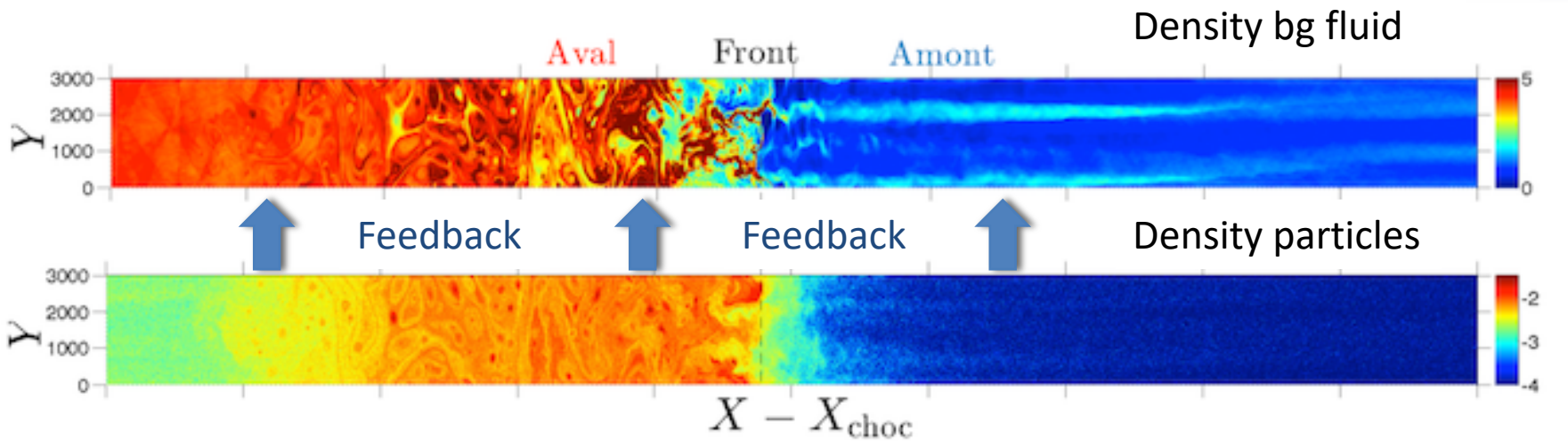
... prompt MeV radiation: dissipation of jet bulk kinetic (magnetic?) energy

... afterglow: dissipation of jet energy through a strong collisionless relativistic shock with the surrounding medium
shock heating of swept up electrons and shock acceleration

Why numerical simulations?

Goal: understand **non-linear** plasma physics of particle acceleration

Beyond standard DSA: non-linearities



If no strong feedback \rightarrow test-particle DSA is valid.

BUT, in general, we do expect strong feedback from accelerated particles...

Different plasma simulation approaches

Types of numerical simulations

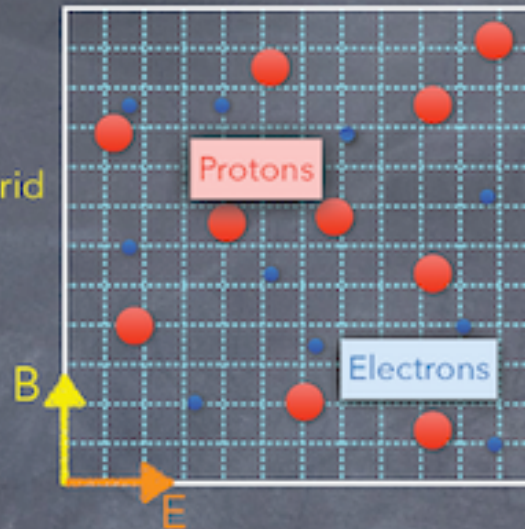
1. Fully kinetic: Particle-In-Cell or Vlasov codes
2. Hybrid-PIC or hybrid-Vlasov: electrons are fluid, ions are macro-particles
3. Hybrid-MHD: either supra-thermals are test-particles (not very interesting...) or coupled through PIC algorithm to the background fluid (MHD+CR as particles from supra-thermals).

Full-PIC and Hybrid-PIC

Full-PIC and Hybrid-PIC techniques

Full-PIC approach

- Define electromagnetic fields on a **grid**
- Move particles via **Lorentz force**
- Evolve fields via **Maxwell equations**
- Computationally very challenging!

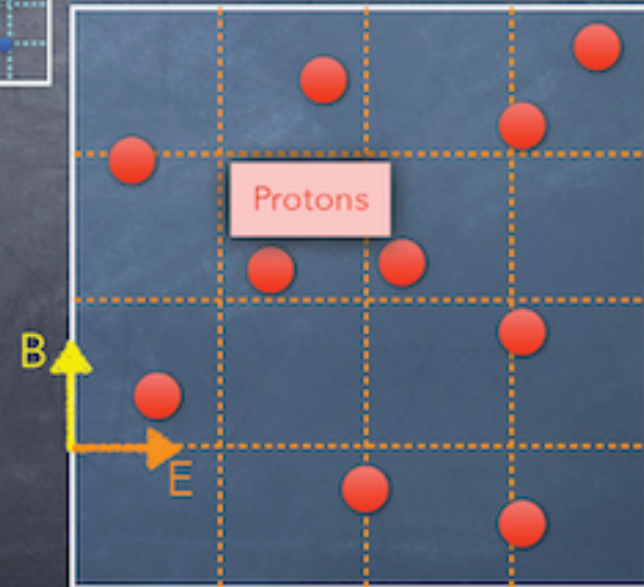


Hybrid approach: Fluid electrons - Kinetic protons

(Winske & Omidi; Burgess et al., Lipatov 2002; Giacalone et al.; DC & Spitkovsky 2013-2018, Haggerty & DC, 2019...)

- massless electrons for more **macroscopical** time/length scales

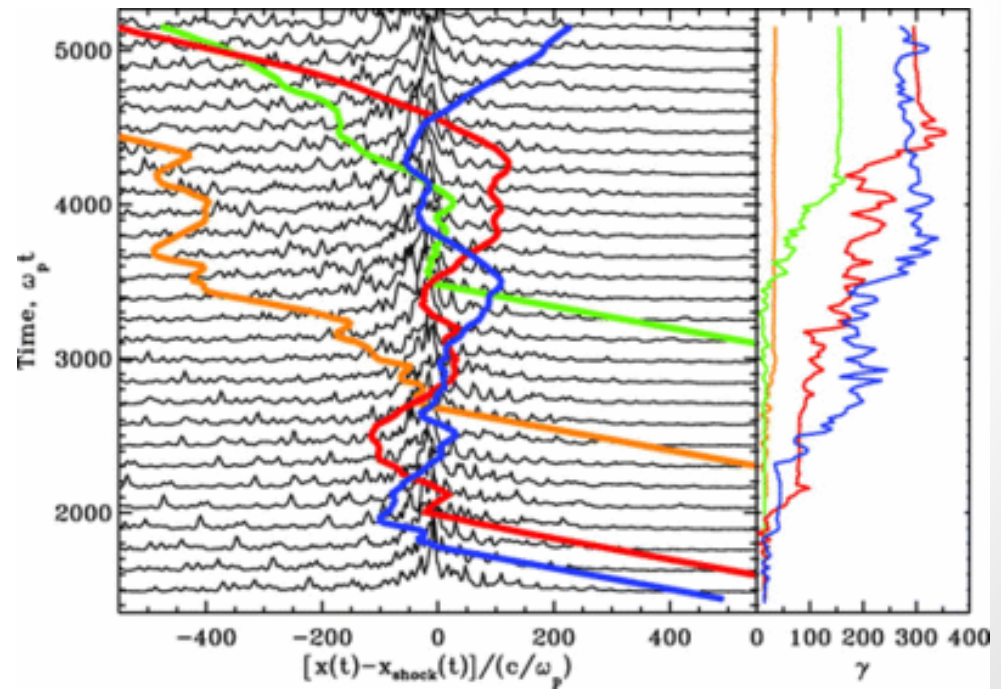
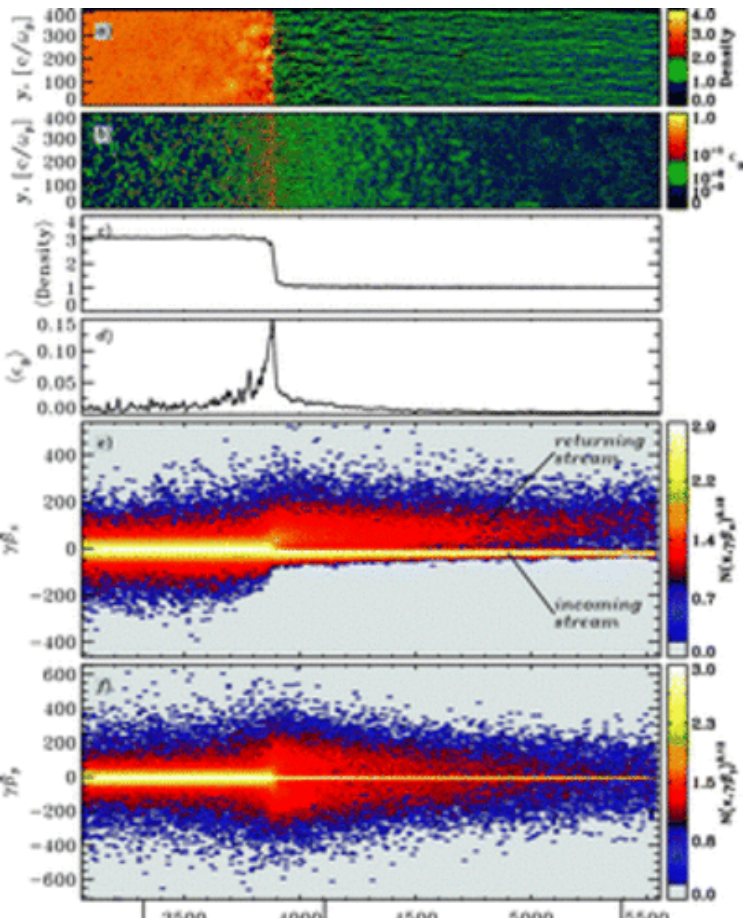
$$\mathbf{E} = -\frac{\mathbf{V}_i}{c} \times \mathbf{B} + \frac{1}{4\pi n e} (\nabla \times \mathbf{B}) \times \mathbf{B} - \frac{T_e}{n} \nabla n^{\gamma_e}$$



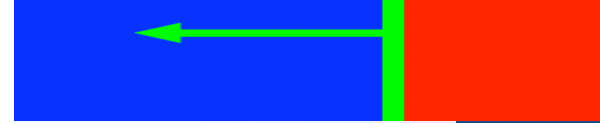
Relativistic shocks simulations

First ab-initio demonstration of Fermi I process

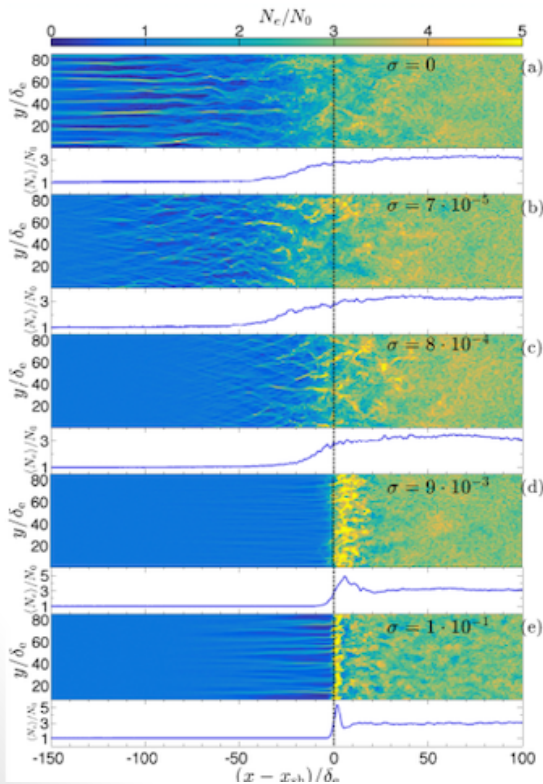
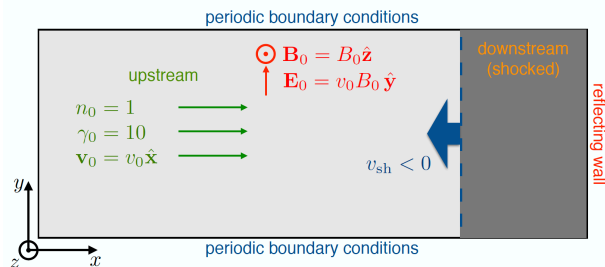
Spitkovsky, ApJL, 2008



Relativistic shocks simulations

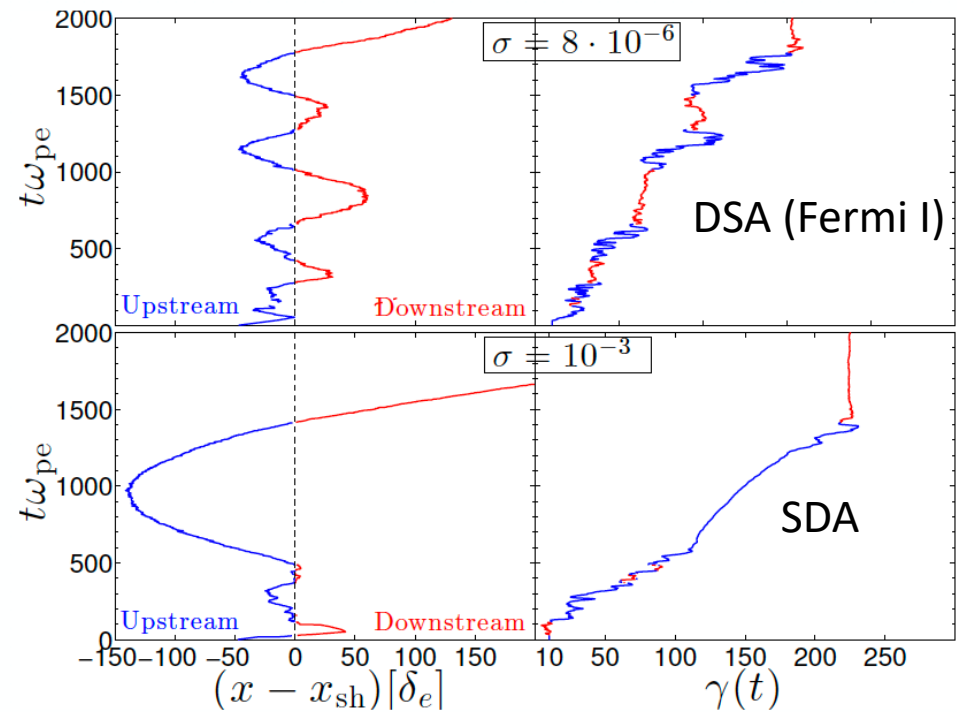


Fermi I process: dependence on magnetization



increasing magnetization

Sironi, Spitkovsky, Arons, ApJ, 2013
 Plotnikov, Grassi, Grech, MNRAS, 2018
 -> Dependence on upstream plasma magnetization

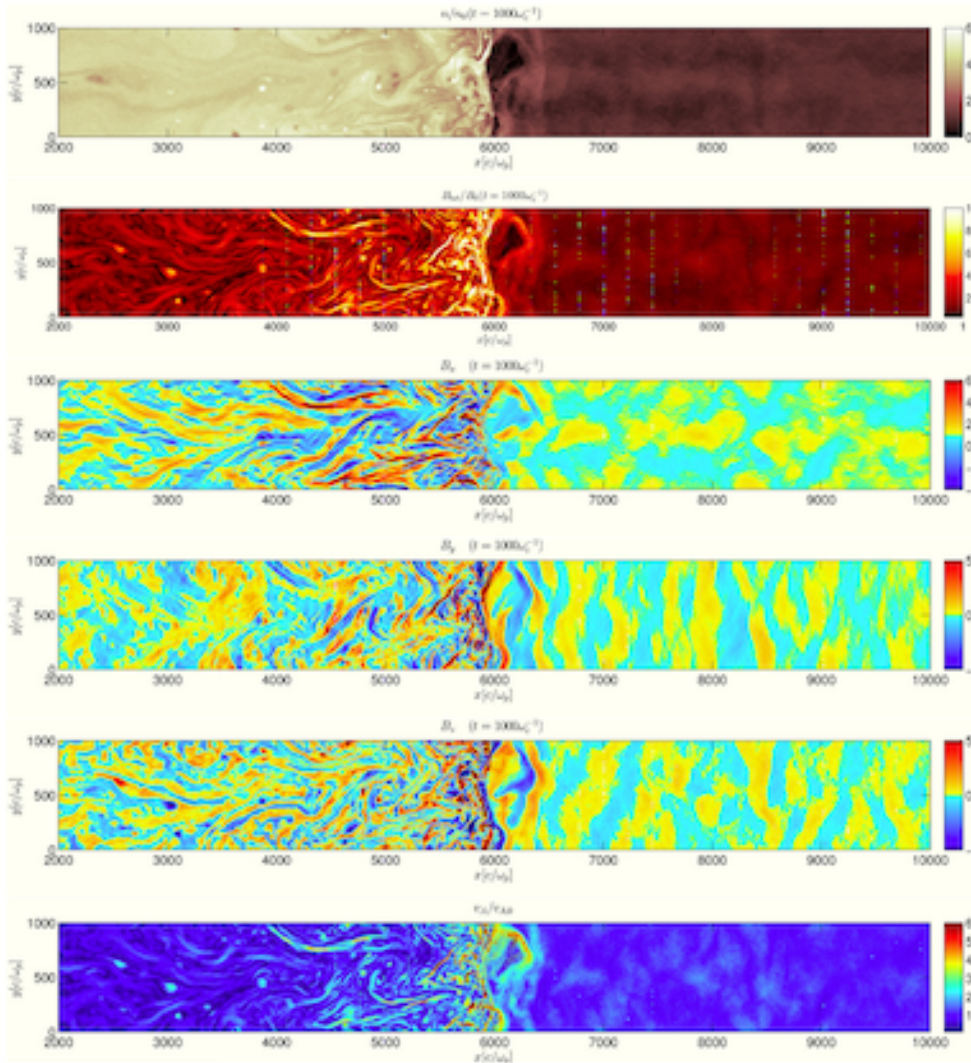


Hybrid-PIC simulation of shock



Global view: $M_A = 20$

Caprioli & Spitkovsky 2014b



D_{sty}

B_{tot}

B_x - 1

B_y

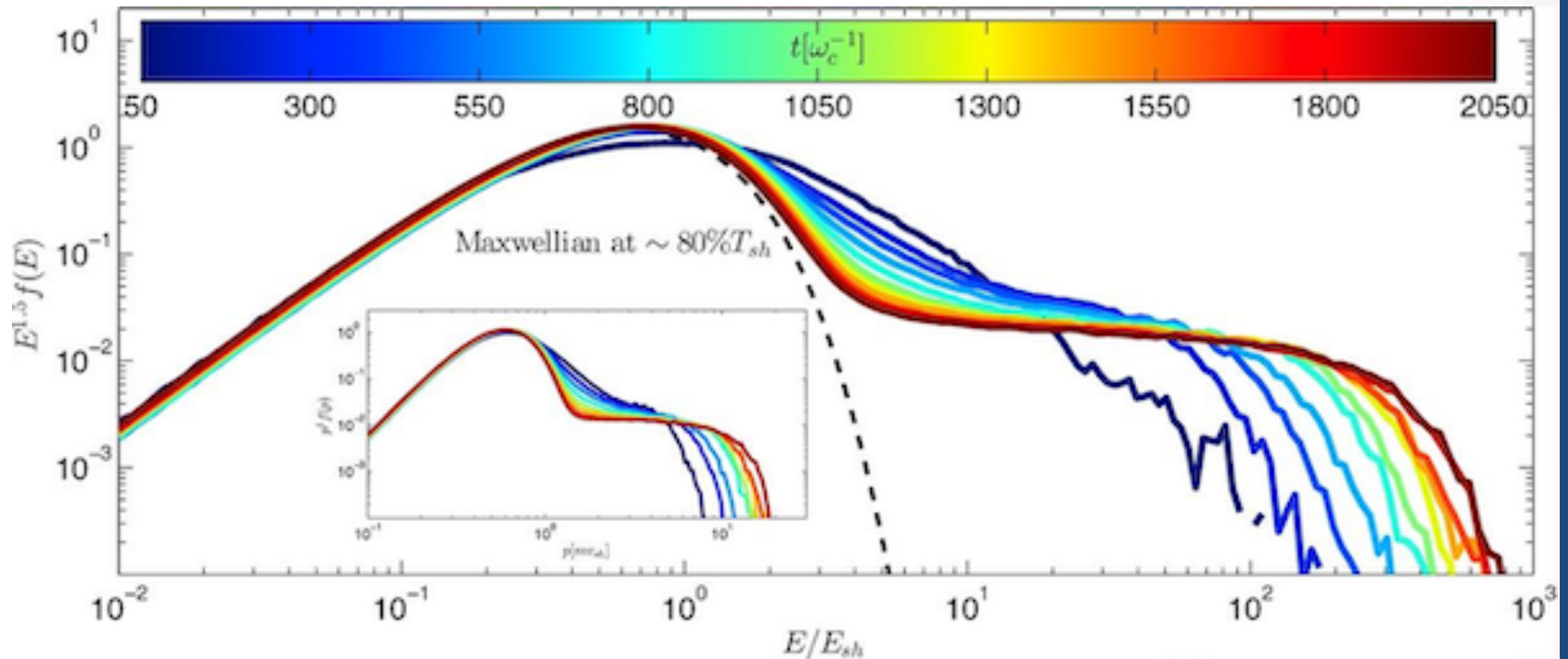
B_z

V_A

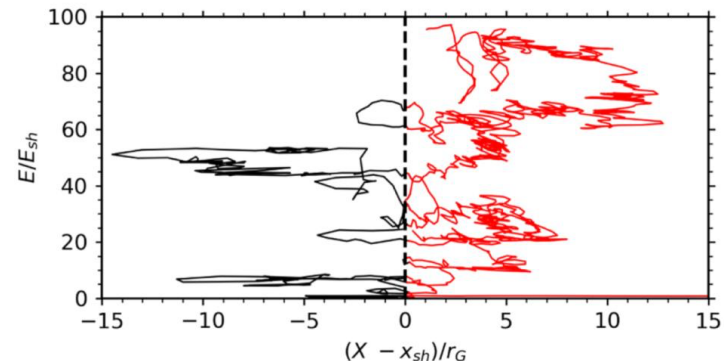
Amplification in the precursor
-> Main non-linearity

Hybrid-PIC simulation of shock

Particle acceleration properties



As particles are accelerated, they escape upstream of the shock.
Leads to efficient B-field amplification.



Caprioli & Spitkovsky 2014a,b,c;
Haggerty et al, 2020; Caprioli et al, 2020

Shocks plasma physics: Not only particle acceleration.

Related question: Magnetic field amplification at shocks

Quasi-perp shocks studied by PIC simulations (Bohdan et al, PRL, 2021)
Evidence of the role of Weibel instability even in non-relativistic shocks.
(Here compared to Saturn's bow shock crossings)

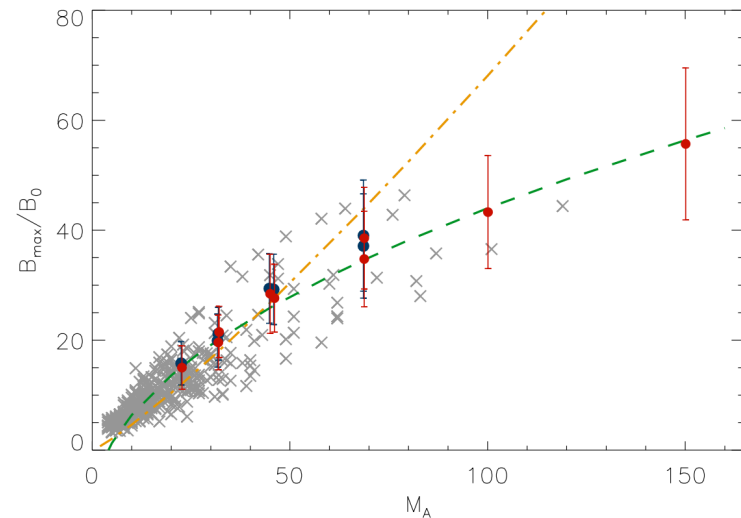
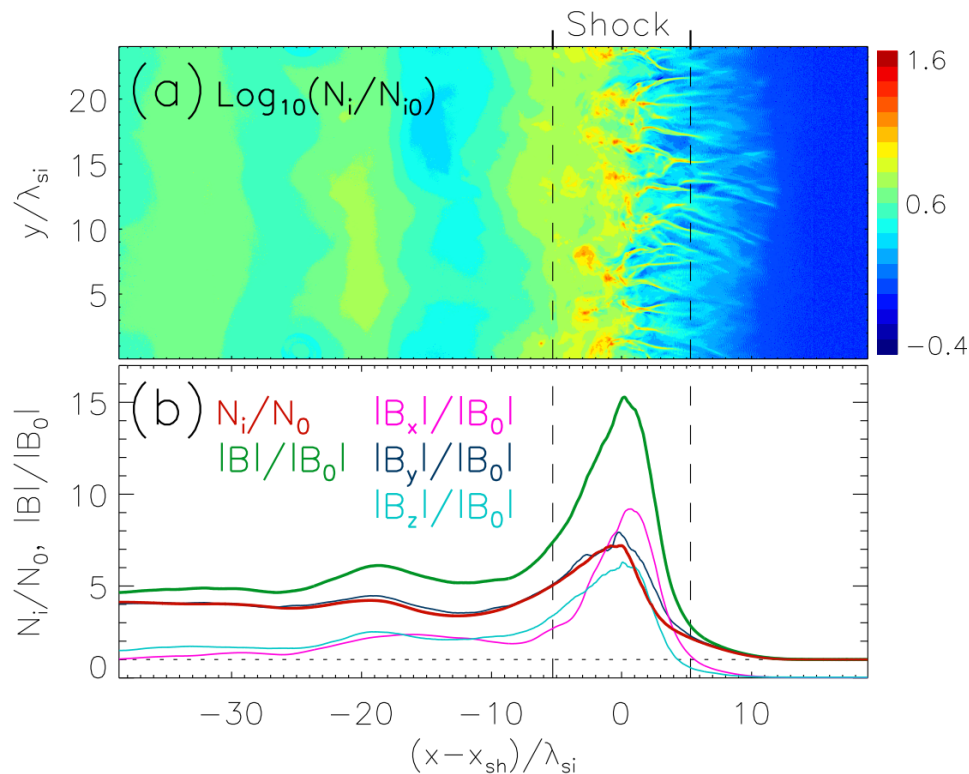


FIG. 4. Cassini measurements [25] indicated by gray crosses and PIC simulation data displayed with blue and red dots for left ($\beta = 5 \cdot 10^{-4}$) and right ($\beta = 0.5$) shocks, respectively. The yellow dash-dotted line is an earlier prediction, $B_{\text{over}}/B_0 \approx 0.4M_A^{7/6}/1.26$ (cf. Eq. 1), corrected for shock reformation. The green dashed line is the behavior found in our PIC simulations, $B_{\max}/B_0 = 5.5(\sqrt{M_A} - 2)$.

MHD-PIC

MHD-PIC summary: MHD with CR particles

Full equations for the CR particles:

$$\frac{d(\gamma_j \mathbf{u}_j)}{dt} = \frac{q_j}{m_j} \left(\mathbf{E} + \frac{\mathbf{u}_j}{c} \times \mathbf{B} \right)$$

Relativistic Boris pusher, subcycling (~ 10 particle steps per MHD).

Specify the numerical speed of light $c \gg$ any velocities in MHD.

Full equations for the gas:

$$\frac{\partial \rho \mathbf{v}}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B} + \mathbf{P}^*) = - \text{Lorentz force on the CRs}$$

$$\frac{\partial E}{\partial t} + \nabla \cdot [(E + P^*) \mathbf{v} - \mathbf{B}(\mathbf{B} \cdot \mathbf{v})] = - \text{energy change rate of the CRs}$$

Momentum and energy source terms reflect Newton's 3rd law.

CR-induced Hall effect

Electrons are force-free: $\mathbf{E} + \frac{\mathbf{v}_e}{c} \times \mathbf{B} = 0$

Decomposition of current density:

$$\frac{c}{4\pi} \nabla \times \mathbf{B} = \mathbf{J}_{\text{tot}} = n_i q_i \mathbf{v}_i - n_e e \mathbf{v}_e + n_{\text{CR}} q_{\text{CR}} \mathbf{u}_{\text{CR}}$$
$$en_e = q_i n_i + q_{\text{CR}} n_{\text{CR}}$$

Generalized Ohm's law:

$$\mathbf{E} = -\frac{\mathbf{v}_i}{c} \times \mathbf{B} + \frac{1}{en_e c} \mathbf{J}_{\text{tot}} \times \mathbf{B} - \frac{q_{\text{CR}} n_{\text{CR}}}{en_e} \frac{(\mathbf{u}_{\text{CR}} - \mathbf{v}_i)}{c} \times \mathbf{B}$$

inductive term

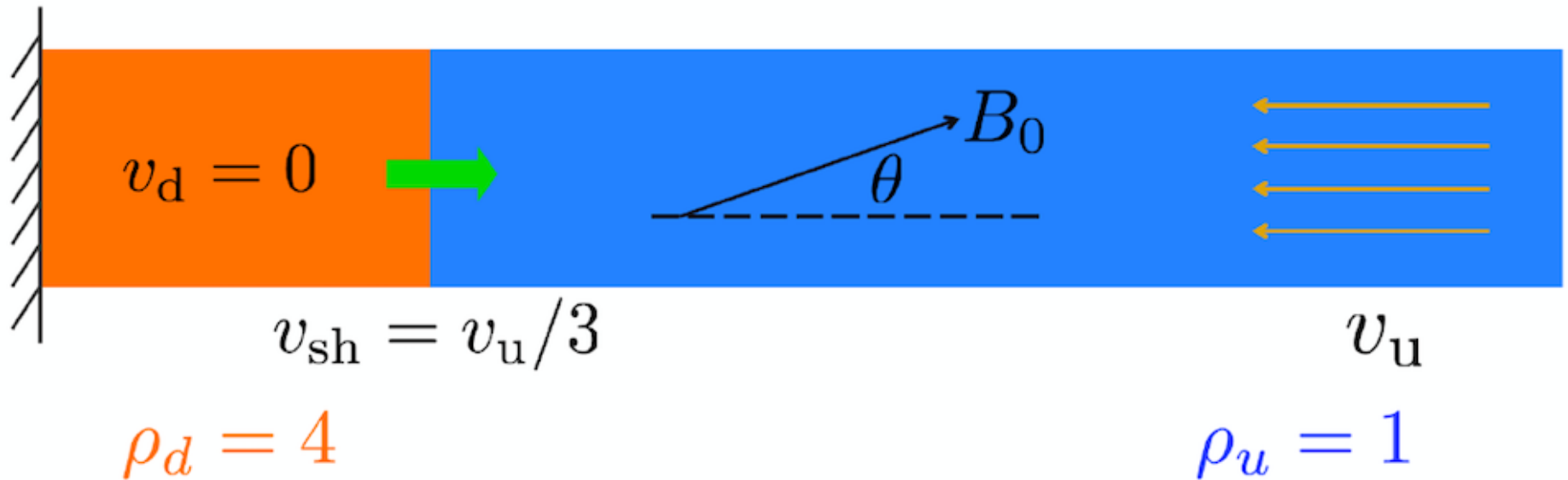
normal Hall term

CR-induced Hall term

Important on scales < ion skin depth

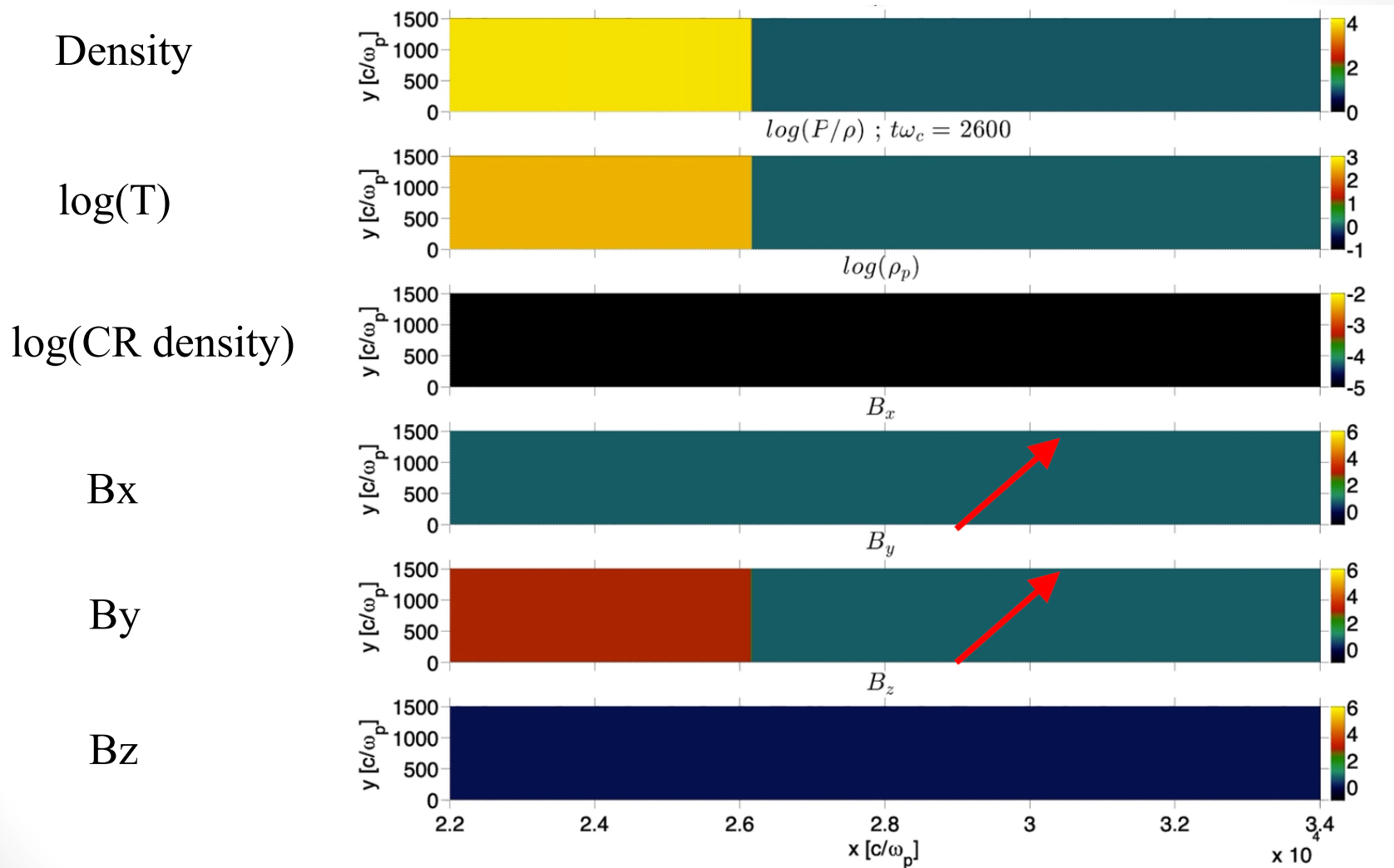
scale independent

Setting up the shock problem



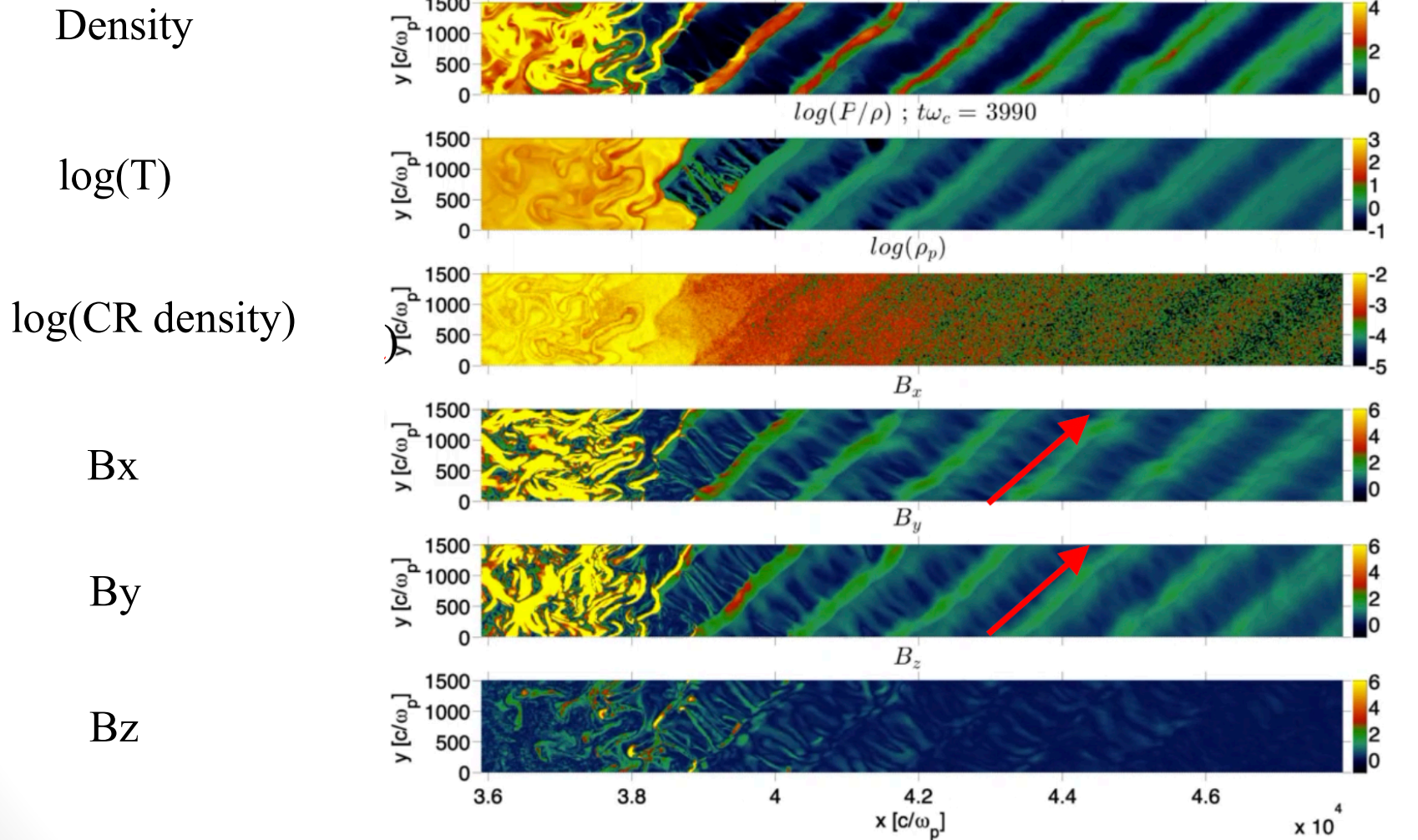
- Inject CR particles at the shock with some efficiency η .
- They are injected at energy of 10 Eshock isotropically.
- Escaping CRs drive upstream waves, and acceleration ensues.

MHD-PIC simulation: oblique shock, no particle injection



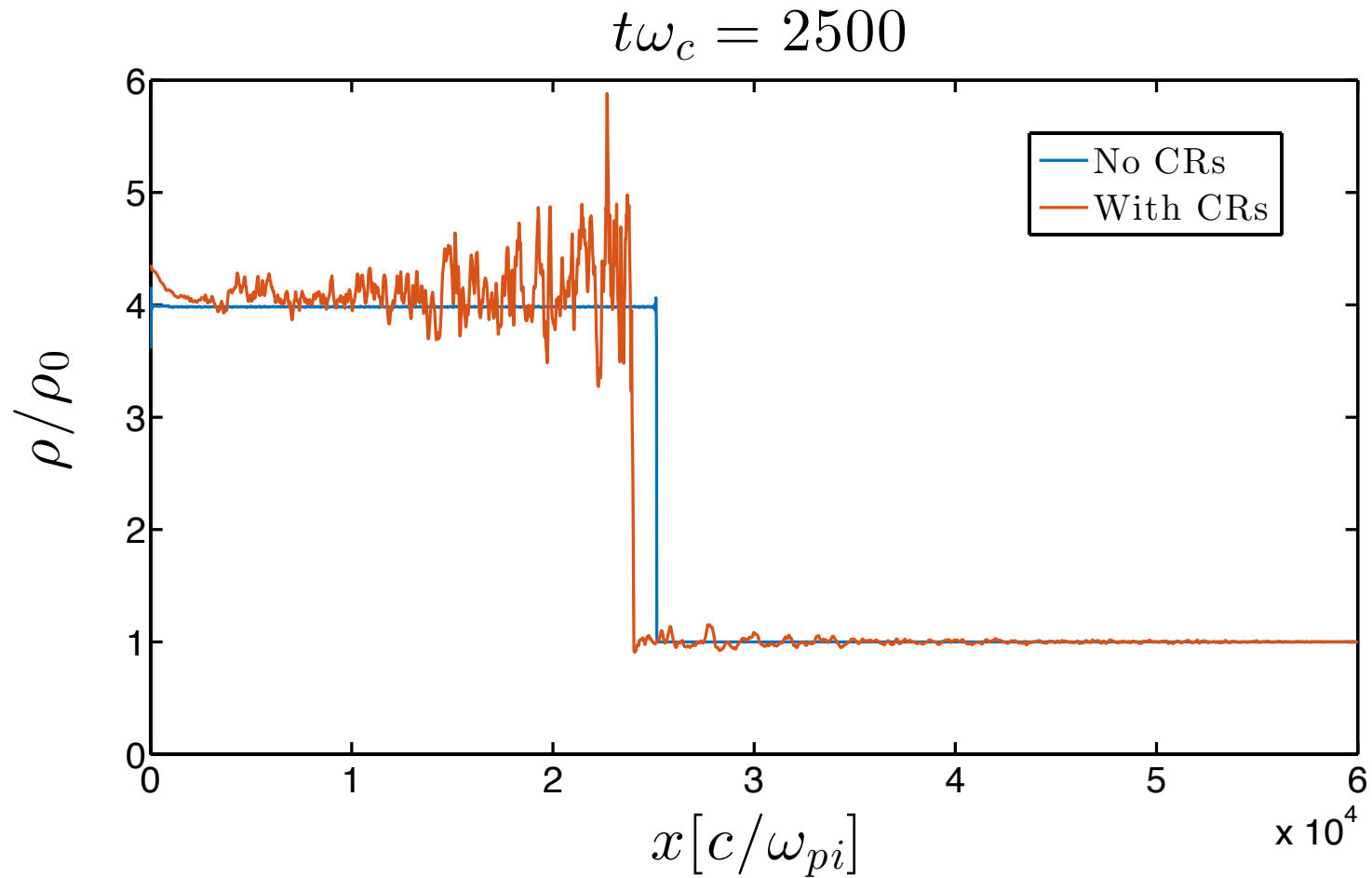
The shock is a jump between two fluid states, consistent with RH conditions

MHD-PIC simulation: injection at front (MA=30 and $\eta_{inj}=2 \cdot 10^{-3}$)



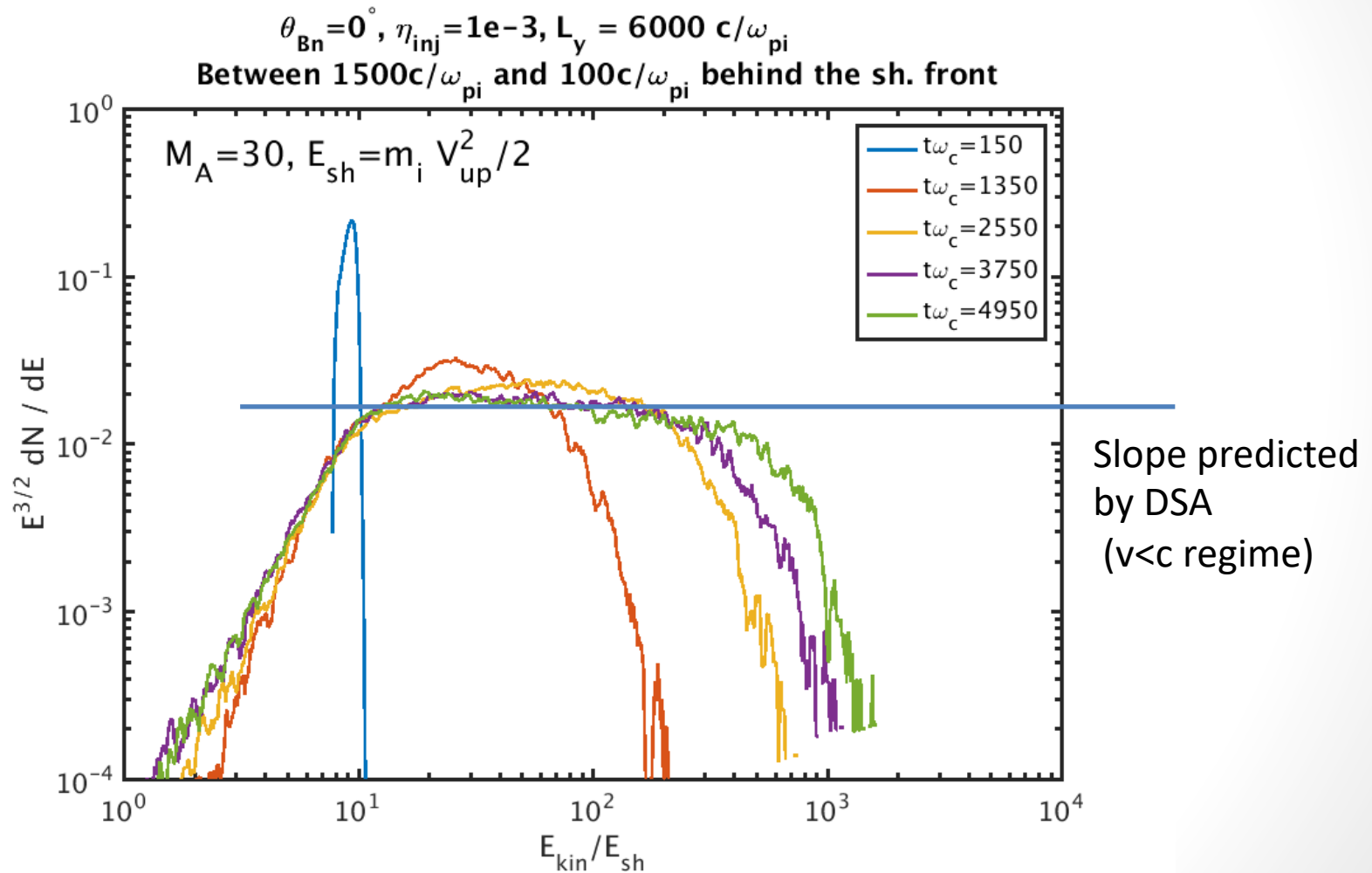
Modified shock: wave amplification upstream, modified jump conditions...

Shock modification by particle acceleration



If efficient acceleration, density compression ratio is enhanced, front slows down

Particle acceleration: power-law build-up

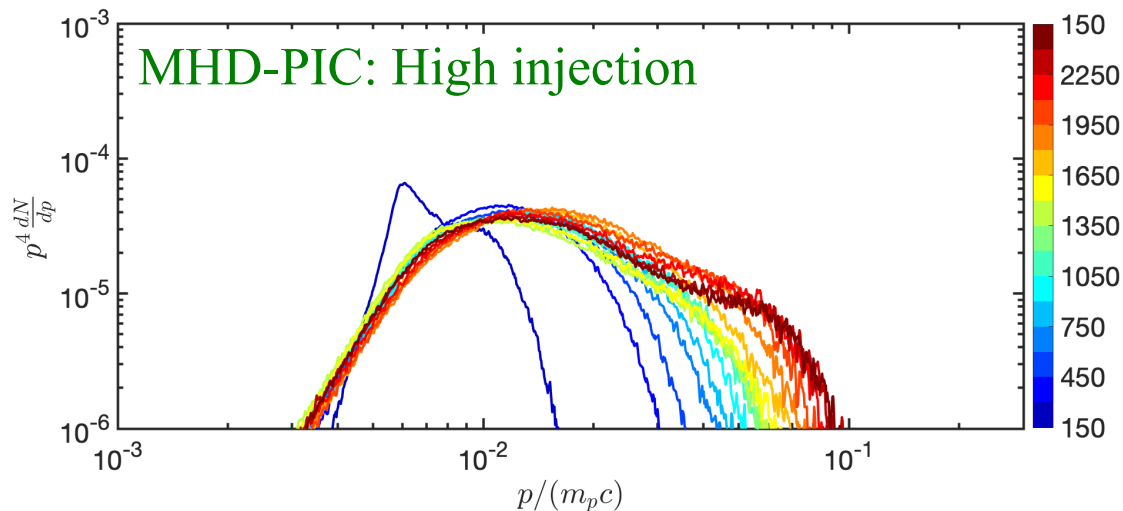
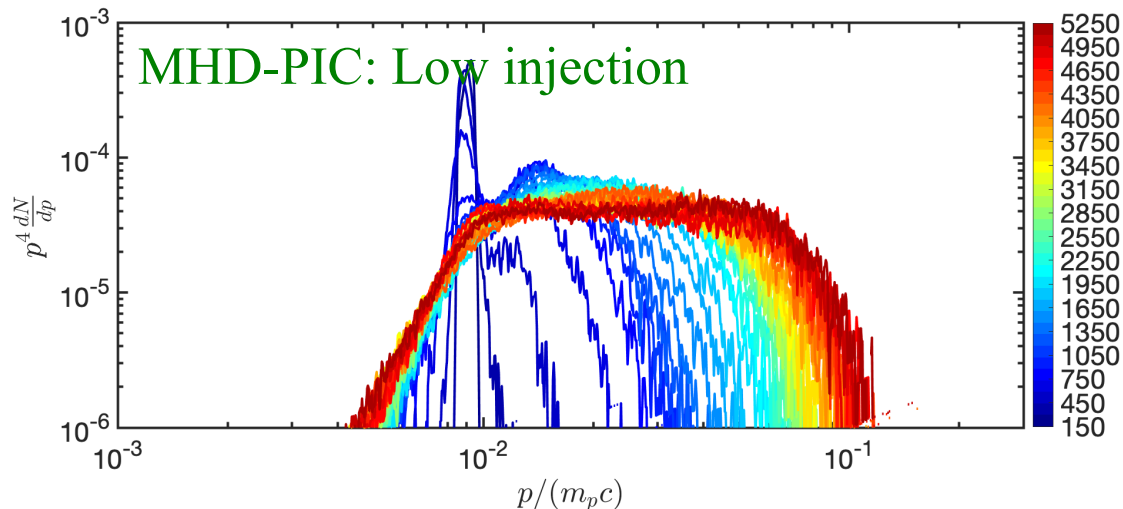


Prominent power-law build up, but no sign of spectrum concavity.

Post-shock $F(p)$: dependence on injection recipe

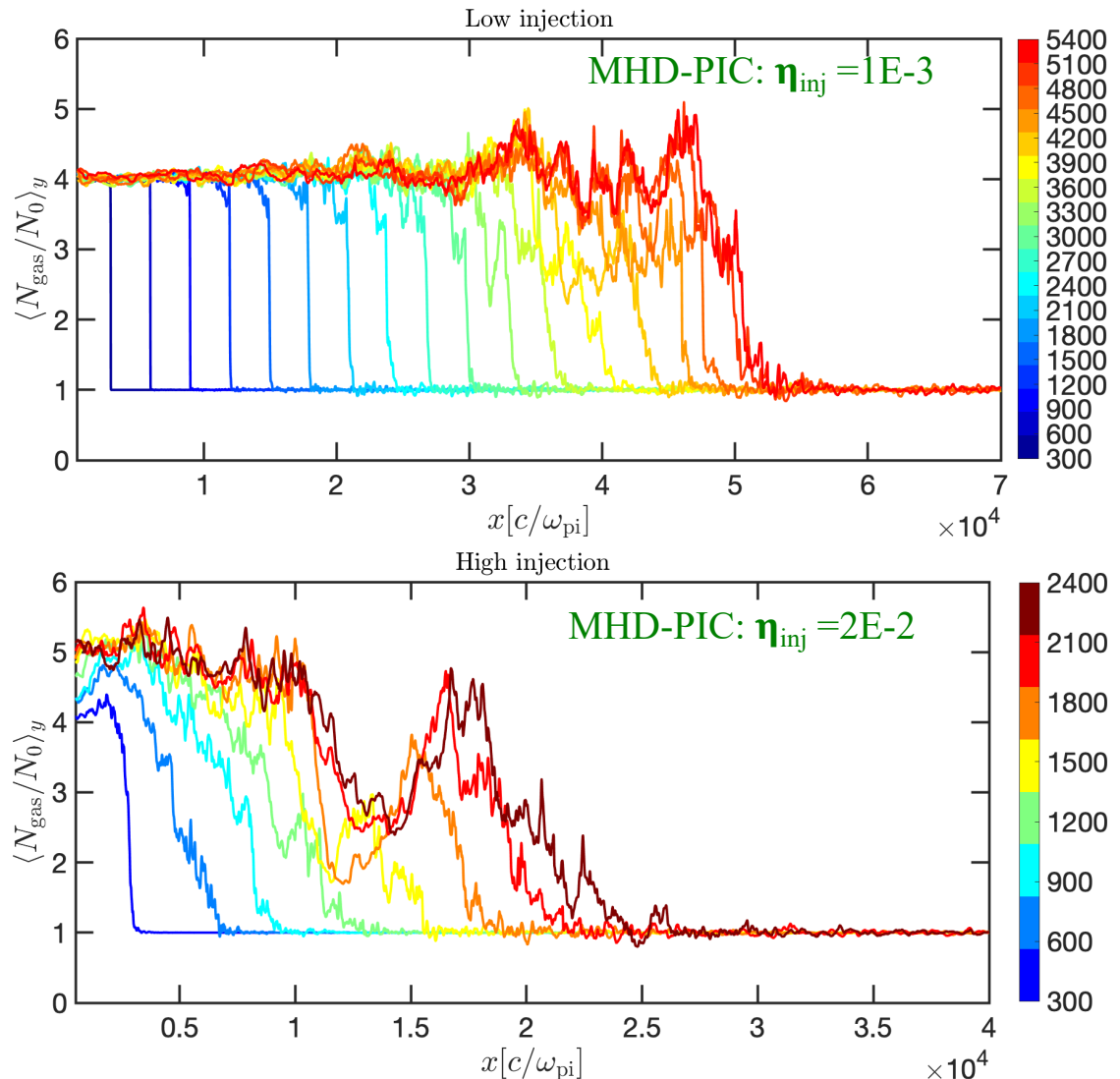
Low injection:
in good agreement with
standard DSA $dN/dp \propto p^{-4}$,
see Caprioli & Spitkovsky
2014a, Bai et al. 2015

Steeper spectrum in the
High injection run.
 $dN/dp \propto p^{-5}$ roughly, at late
times.



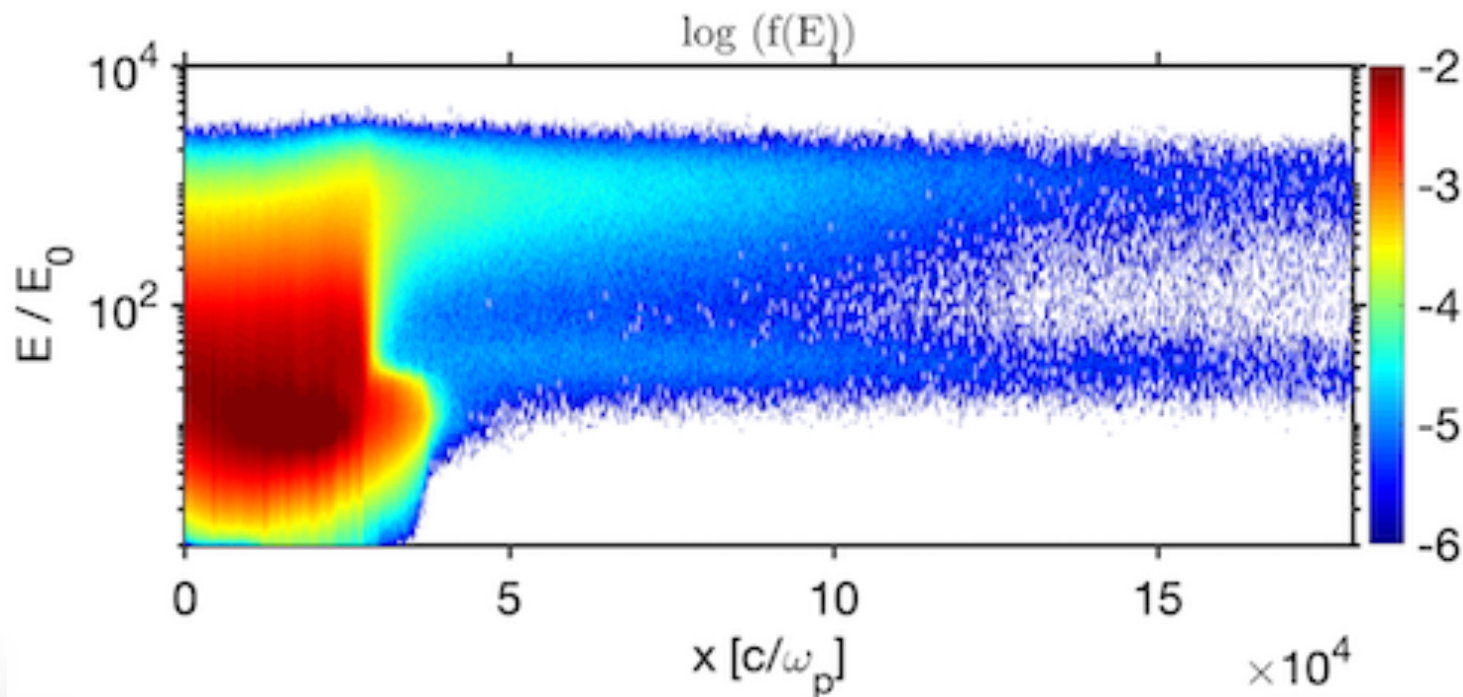
Compression ratio: MHD-PIC runs

- Y-averaged Gas number density at different simulation times (see colorbar in units of cyclotron)
- **Low injection** here ($\eta_{inj} = 1e-3$): Far downstream : ratio 4 conserved as in standard picture. Close downstream: starts to increase at late times ($t > 3000$)
- **High injection** (bottom panel, $\eta_{inj} = 2e-2$). Far downstream : increase to 5. Close downstream : strong modification leads to poor definition of the shock front



Phase space evolution of CRs

- Injected behind the front with $E = 10 \cdot E_0$
- Movie below
- CRs escape upstream and produce waves
- Self-confinement at late times



Summary and open questions

Micro and multi-scale simulations of particle acceleration

- Micro scale: powerful ab-initio plasma simulation techniques. Important to study initial stages but fail to bridge the gap with astrophysical scales.
- MHD-PIC: aims to bridge the gap in almost self-consistent way. The difficulty consists in a few prescriptions.
- A lot of open questions to be investigated.

Further reading:

- Marcowith et al, 2020, LRCA, 6, 1
« Multi-scale simulations of particle acceleration in astrophysical systems »



Review on multi-scale simulations

Some links to public PIC codes

PIC codes

- **TRISTAN-MP** (Princeton group)
<https://ntoles.github.io/tristan-mp-pitp/>
- **SMILEI** (French, Polytechnique + Maison de la Simulation)
<https://smileipic.github.io/Smilei/>
- **ZELTRON** (B. Cerutti)
<https://ipag.osug.fr/~ceruttbe/Zeltron/>
- **EPOCH** (UK)
<https://github.com/Warwick-Plasma/epoch>
- **OSIRIS** (not open-source...)
<https://picksc.idre.ucla.edu/software/software-production-codes/osiris/>