Wave Chaos in Rapidly Rotating Stars

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Universality in chaotic systems

The shape of the spacing distribution is roughly the same for all chaotic spectra.



Effect of Rotation

At increased rotation we see new types of p-modes [1]:

- whisperring gallery modes
- 2-period island modes
- 6-period island modes
- chaotic modes

Work has been done to understand

Quantum Chaos

Quantum chaos is the study of quantum systems whose classical limit is chaotic.



Action Along Ray Paths

The results of quantum chaos remain true outside of quantum physics, with :

 $\hbar \to \omega^{-1}$ $S \to T$, where T is the travel time of a ray along a certain path.

Now we want $\rho(T)$.

To find chaotic paths, we look at the



the asymptotic structure of the ilsand modes spectrum [2]. The goal of my thesis is to study the spectrum of chaotic modes.

stationnary wavefunction inside a cardioid cavity

chaotic wavefunctions are ergodic : the probability density is rather uniformly distributed inside the cavity.

Poincaré surface of section. For long times, a single chaotic trajectory will fill up a region of phase space.



Poincaré surface of section : chaotic zone

Keeping only periodic orbits, we end up with the following distribution $\rho(T)$:



Asteroseismology

By studying the oscillations of a star we can gain information about its interior : that is the goal of asteroseismology. In this thesis I will focus on oscillations induced by pressure/acoustic waves (pmodes).



Regularities in the chaotic spectrum

The autocorrelation of the chaotic spectrum shows a big spike at the great separation. It is an unusual feature for a chaotic system.



Chaotic spectrum



Autocorrelation of the chaotic spectrum

oscillation mode of a non rotating star

In the case of slowly rotating stars, the spectrum of acoustic oscillations is well discribed in the high-frequency regime by Tassoul's asymptotic formula :

$$\omega_{n,l} = \Delta \omega \left(n + \frac{\ell}{2} + \frac{1}{4} + \alpha \right)$$
$$\Delta \omega = 2\pi \left(2 \int_0^{R_e} \frac{dr}{c_s} \right)^{-1}$$

This formula relates the great

Ray Dynamics

In the small wavelength approximation acoustic waves propagate according to the eikonal equation :

 $\omega^2 = \omega_c^2 + c_s^2 k^2$

From this we derive the Hamiltonian :

 $H = \frac{\mathbf{k}^2}{2\omega^2} + W$ $W = -\frac{1}{2c_c^2} \left[1 - \frac{\omega_c^2}{\omega^2} \right]$

The potential barrier at the surface, responsible for the back-reflection of the wave, stems from the sharp increase of the cutoff frequency ω_c .

An important result of quantum chaos, that is valid for bounded systems, is the Gutzwiller formula :

$$d(E) = \overline{d}(E) + \frac{1}{\hbar} \sum_{j} A_{j} e^{i\left(\frac{S_{j}(E)}{\hbar} + \nu_{j}\right)}$$

The sum is calculated over the periodic orbits : S_i is the action along the jth orbit and A_i is linked to the stability of the jth orbit.

Having an expression for the density of the spectrum might seem enough to calculate other statistical quantities, such as variance, but the series is divergent and therefore hard to

Distribution of the actions along periodic orbits

Conjecture : the oscillation of $\rho(T)$ is at the origin of the chaotic spectrum correlations.

References

[1] F. Lignières, B. Georgeot. Asymptotic analysis of high-frequency acoustic modes in rapidly rotating stars A&A 2009

[2] M. Pasek, F. Lignières, B. Georgeot, D. R. Reese.

Regular oscillation sub-spectrum of rapidly rotating stars

separation $\Delta \omega$, a quantity we can access through observations, to the speed of sound c_s in the star.



Solar oscillations





Chaotic trajectory of an acoustic ray

manipulate.

Michael Berry showed [3] that the fourier transform K(T) of the correlation function $< \left[d(e - \frac{1}{2}\xi) - 1\right] \left[d(e + \frac{1}{2}\xi) - 1\right] > can$ be approximated as :



Where ρ is the distribution of actions along periodic orbits.

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[3] M. Berry Some quantum to classical asymptotics

