

Main sequence evolution of rapidly rotating early-type stars

Journée des thèses

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- Rotation...
 - breaks the 1D spherical symmetry
 - is responsible for large scale mixing processes of chemical elements and angular momentum
- Typical 1D models : rotation \equiv small perturbation
 - Justified for slow rotators (Zahn 1992)
- Early-type stars : often fast rotators
 - 2D evolution possible with the ESTER code (Espinosa Lara & Rieutord 2013; Rieutord et al. 2016)

- Massive stars ($M > 7M_{\odot}$) have winds driven by radiation
- Continuum & line absorptions transfer radiation momentum to stellar matter
 - Matter acceleration outward
 - Mass & angular momentum loss
- At each step, we decrease the mass and the global angular momentum of the star

- Approximations (CAK 1975)
 - Radiative wind : isothermal & stationary flow
 - Star : point source of radiation
- Derivation based on Bard & Townsend 2016 method
- Only need mass and momentum conservation equations

After a mathematical derivation for $c_s \ll 1$:

$$\dot{m}(\theta) = \frac{\alpha}{1 - \alpha} \left(\frac{\bar{Q}\Gamma_e(\theta)}{1 - \Gamma_e(\theta)} \right)^{\frac{1-\alpha}{\alpha}} \frac{\sigma T_{\text{eff}}(\theta)^4}{c^2} \Phi_R(\theta)$$

where

$$\Phi_R(\theta) = \left(1 - \frac{R(\theta)^3 \Omega(\theta)^2 \sin \theta}{GM(1 - \Gamma_e(\theta))} \right)^{-\frac{1-\alpha}{\alpha}}$$

is the rotation term and

α is the CAK index that we take as constant : $\alpha = 0.5$

\bar{Q} is Gayley 1995 dimensionless line strength parameter ($\sim 2 \times 10^3$)

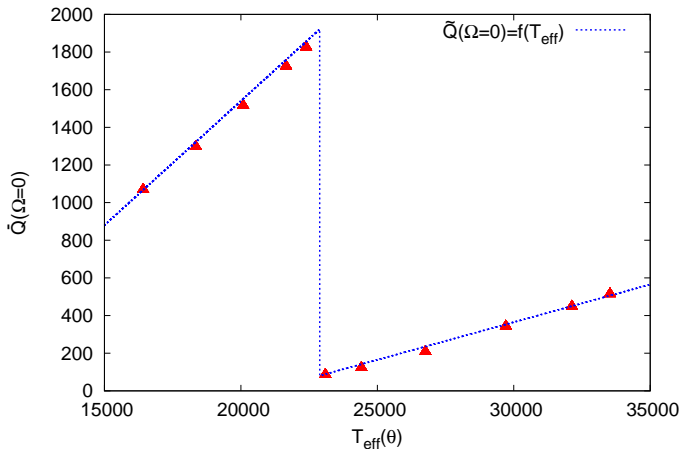
\bar{Q} calibration and bi-stability jumps

- Vink et al. 1999 : recombination of Fe IV into Fe III around $T_{\text{eff}}^{\text{jump}} \simeq 25\,000\text{K} \rightarrow \dot{M}$ increases by a factor $\sim 5 - 10$
- $\bar{Q} = \text{cste}$ ignores this phenomenon
- \bar{Q} has to take the radiative acceleration jump into consideration around this temperature

→ \bar{Q} is calibrated with Vink et al. 2001 prescription for mass-loss in the non-rotating limit ($\Phi_R = 1$)

$$\frac{\dot{M}_{\text{Vink}}}{4\pi R(\theta)^2} = \dot{m}(\theta, \Phi_R = 1)$$

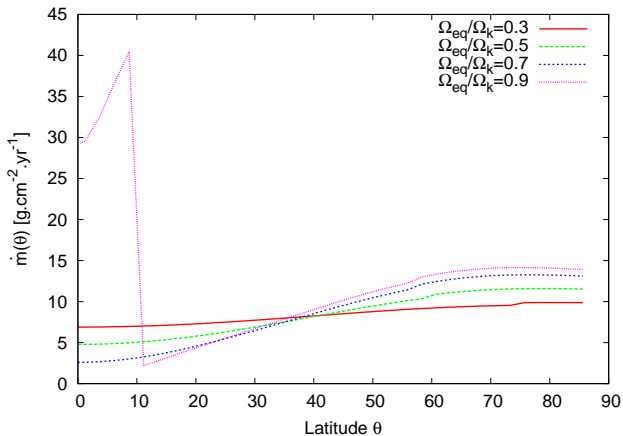
\bar{Q} calibration and bi-stability jumps



$$\tilde{Q}(T_{\text{eff}}) \simeq \begin{cases} 0.132 \times T_{\text{eff}} - 1100, & \text{if } T_{\text{eff}} \leq T_{\text{eff}}^{\text{jump}} \\ 0.04 \times T_{\text{eff}} - 835, & \text{if } T_{\text{eff}} > T_{\text{eff}}^{\text{jump}} \end{cases}$$

Local mass-flux

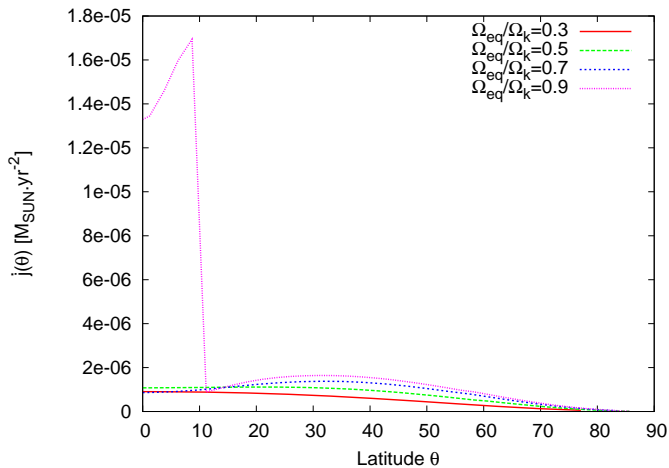
$$M = 15M_{\odot}$$



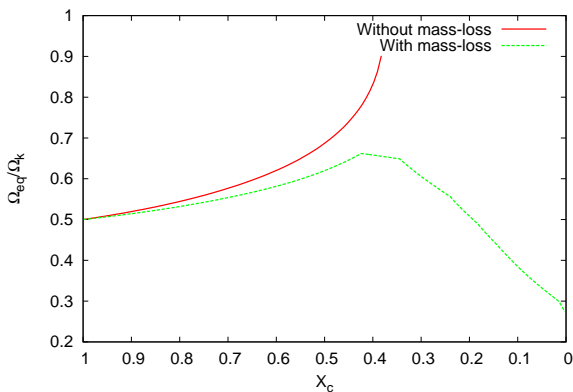
Strong mass-flux in the equatorial region \rightarrow Be stars?

Local angular-momentum flux

$$\dot{j}(\theta) = \dot{m}(\theta)\Omega_{eq}(\theta)R^2(\theta)\sin^2\theta$$



$$\Omega_k = \sqrt{GM/R_{eq}^3}, M = 15M_{\odot}$$



Bi-stability jumps are of crucial interest regarding stellar evolution on the MS

- New mass-flux prescription accounting for bi-stability jumps via \tilde{Q}
- Rotation induced bi-stability :
 - Enhanced mass-flux for $\theta < \theta_{\text{jump}}$
→ Be stars disk ?
 - Enhanced global loss of angular momentum
→ $\frac{\Omega_{eq}}{\Omega_k}$ drop

→ 2D models essential for rapidly rotating stars evolution